

Experimental methods in particle physics

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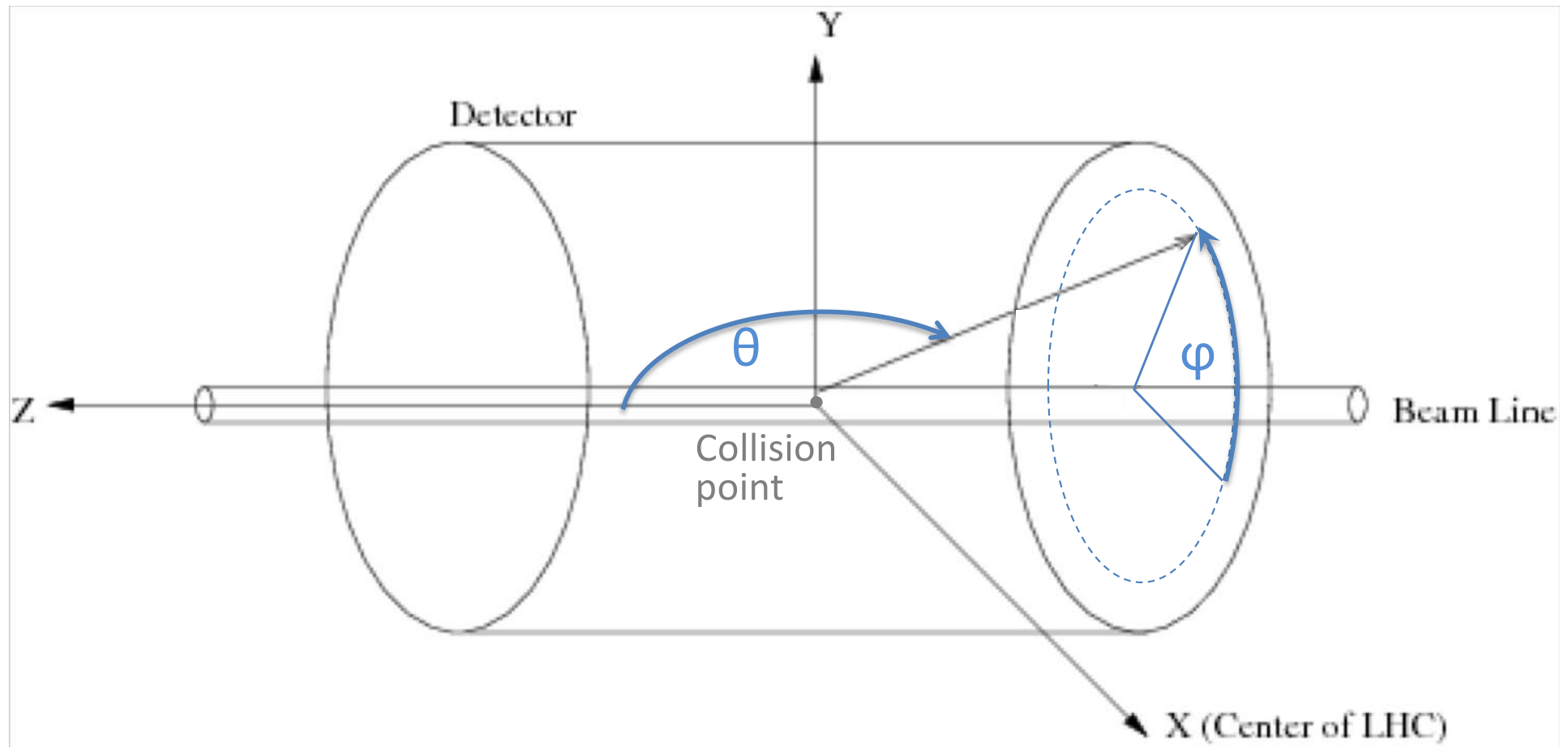
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Webpage of lecture:

<http://atomfizika.elte.hu/rfkm/rfkm2019.html>

Coordinate system of large detectors at colliders

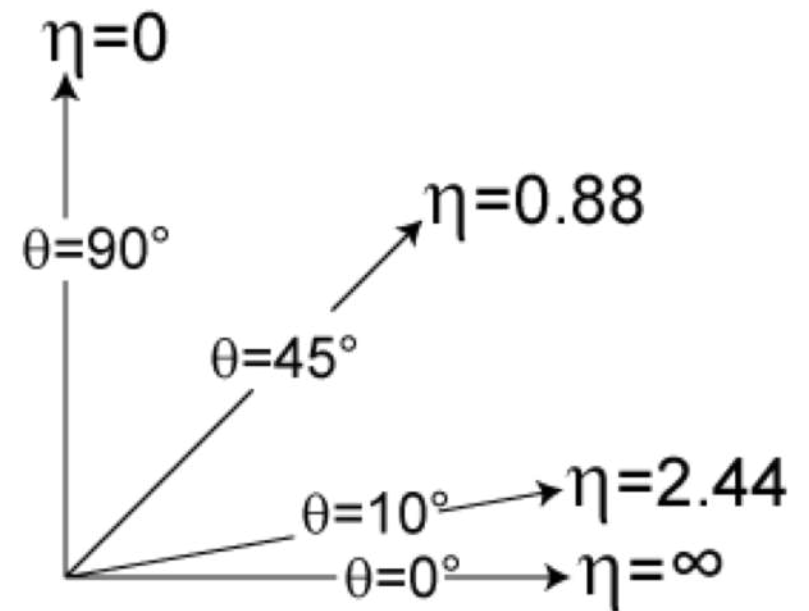
$$\eta > 0 \longleftrightarrow \eta < 0$$



pseudorapidity: $\eta = -\ln \tan (\theta/2)$

Pseudorapidity (η)

- Spatial coordinate
- Angle measured from beam line
- Θ : \angle (particle \mathbf{p} , beam \mathbf{z})
- p_L : longitudinal momentum (p_z)
- If $v \rightarrow c$, $\eta \rightarrow y$ (rapidity)
- Particle production distribution roughly constant as a function of η
- (Pseudo)rapidity difference Lorentz invariant
- ΔR : Lorentz-invariant quantity to measure the angular distance of two particles
- ϕ : azimuth angle
- “Forward” direction: large η , close to beam direction
- “Backward” $-z$ direction (when distinguished from “forward” $+z$ direction)



$$\eta \equiv -\ln \left[\tan \left(\frac{\theta}{2} \right) \right]$$

$$\theta = 2 \arctan \left(e^{-\eta} \right)$$

$$\eta = \frac{1}{2} \ln \left(\frac{|\mathbf{P}| + p_L}{|\mathbf{P}| - p_L} \right) = \operatorname{artanh} \left(\frac{p_L}{|\mathbf{P}|} \right)$$

$$y \equiv \frac{1}{2} \ln \left(\frac{E + p_L}{E - p_L} \right)$$

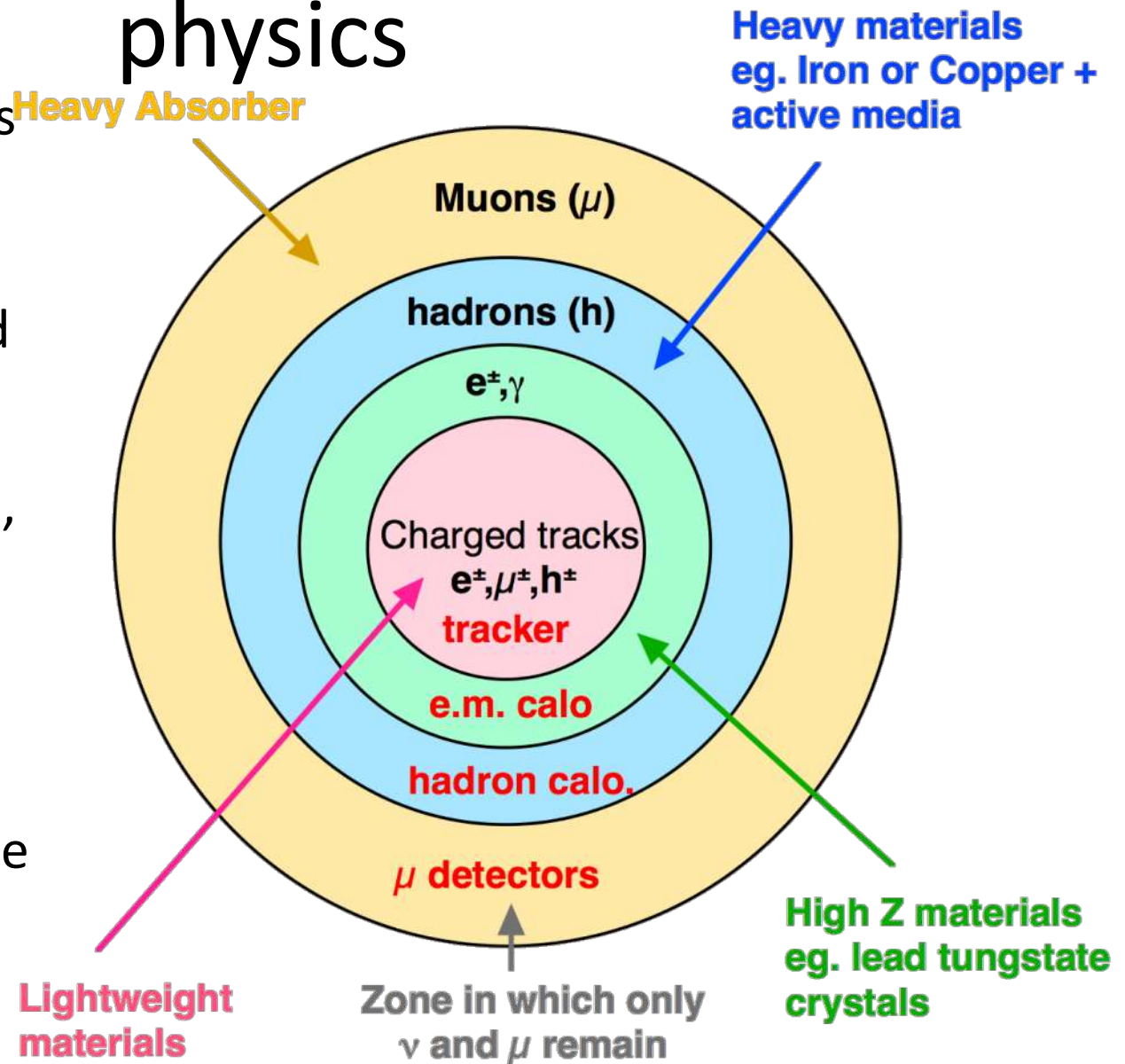
$$m \ll p \Rightarrow E \approx p \Rightarrow \eta \approx y$$

$$(\Delta R)^2 \equiv (\Delta \eta)^2 + (\Delta \phi)^2$$

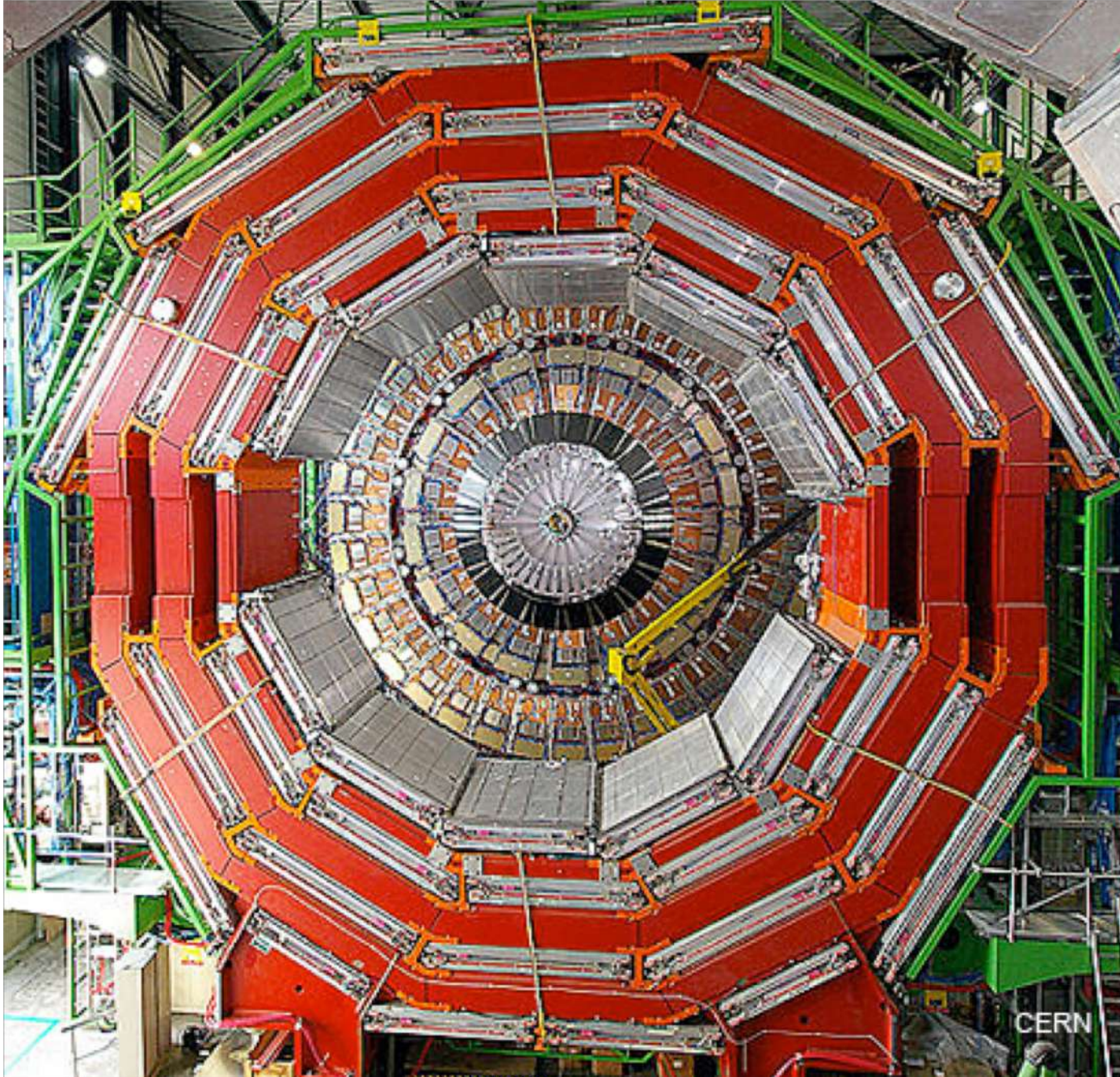
Detector systems in high energy

physics

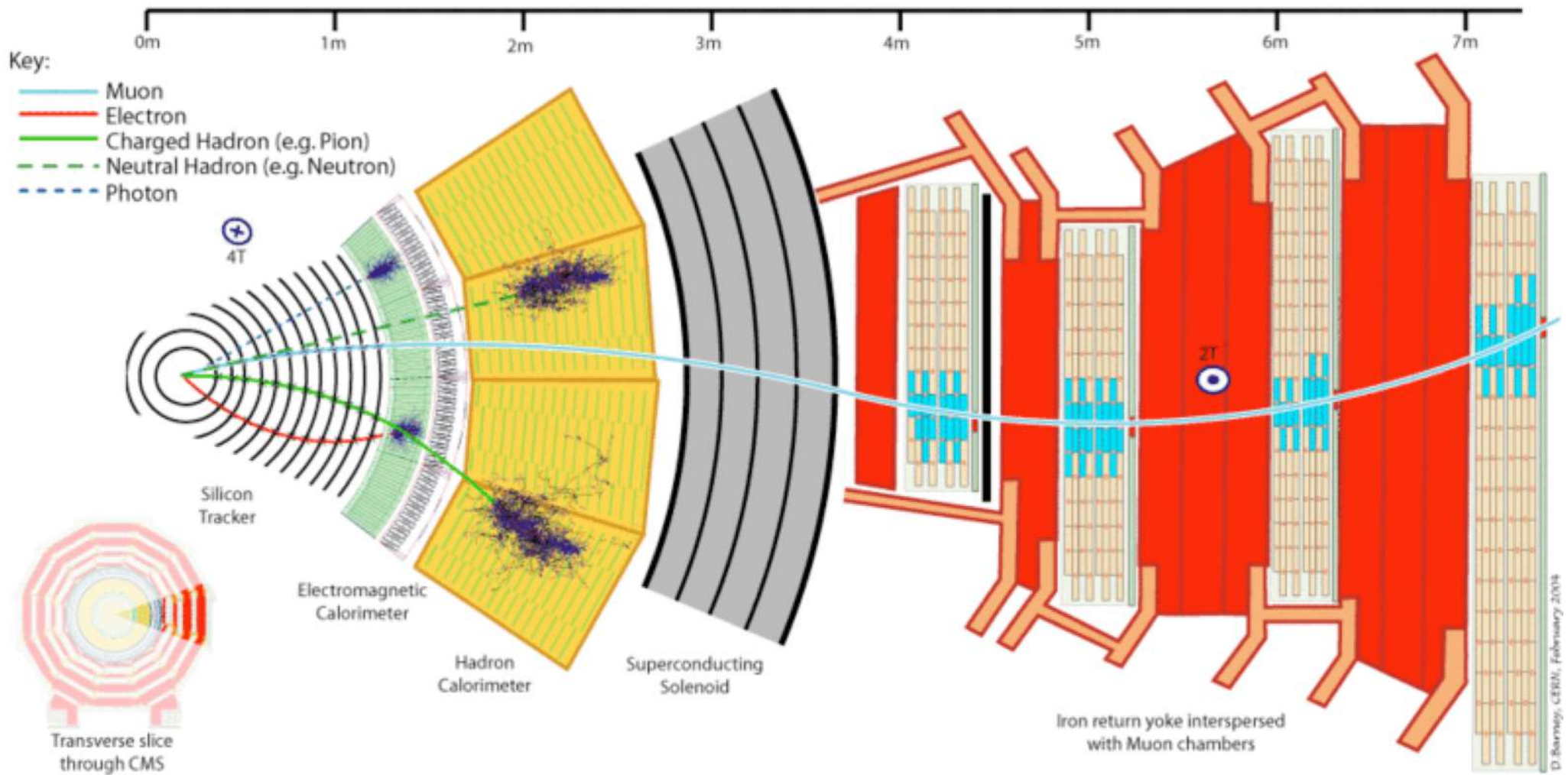
- Detector: an apparatus that measures a particle's position, arrival time, momentum, energy, properties (m, lifetime, quantum numbers...) and identifies it
- There is no detector that could measure / identify, all kind of particles
- Onion-style structure: each layer brings new information
- Comparing the signals of the various detectors, one can identify the particles and measure their momentum, energy



CMS detector



CMS detector

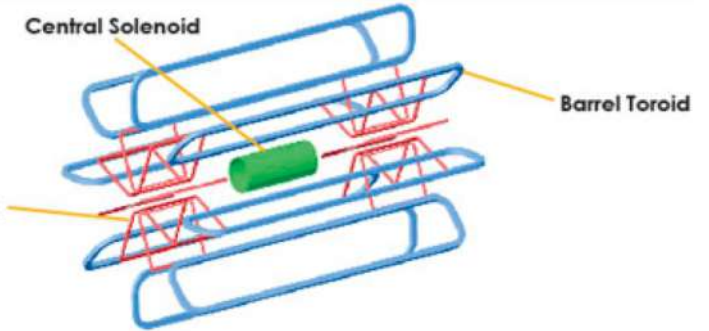


Comparing the signals left by the particles in different detector layers, we can identify the particles and measure their momentum / energy

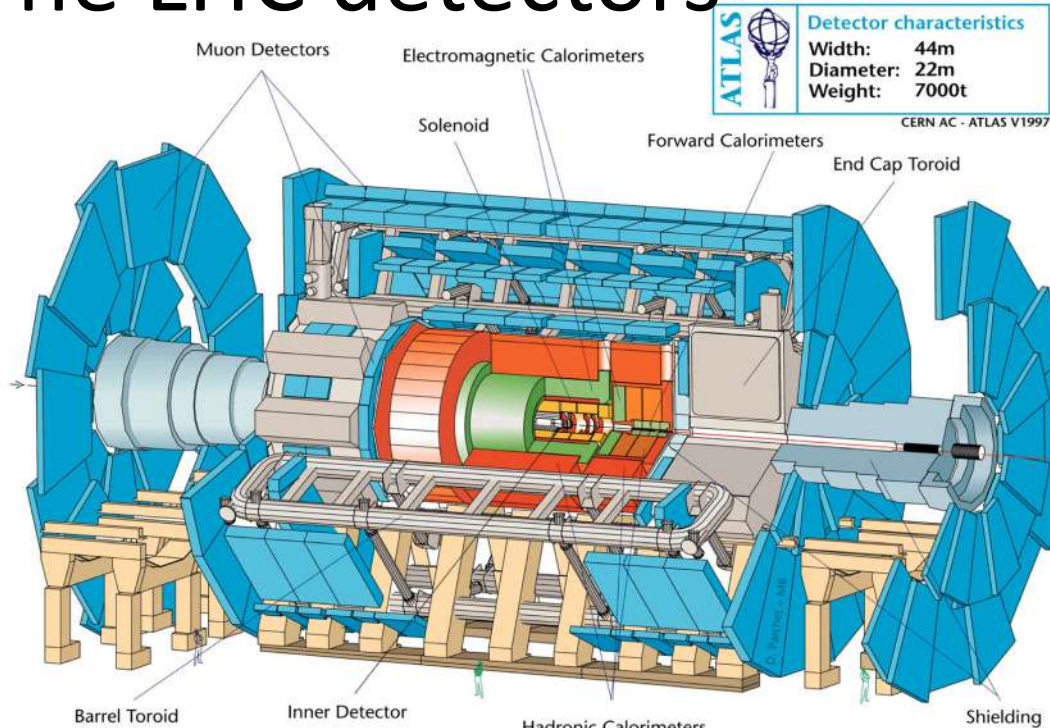
The LHC detectors

Solenoid: 2T
Toroid: 4T on superconductor
non-uniform field

ATLAS



Different design philosophy
Similar resolution, efficiency



Detector characteristics	
Width:	44m
Diameter:	22m
Weight:	7000t

CERN AC - ATLAS V1997

~100 millió elektronikus csatorna

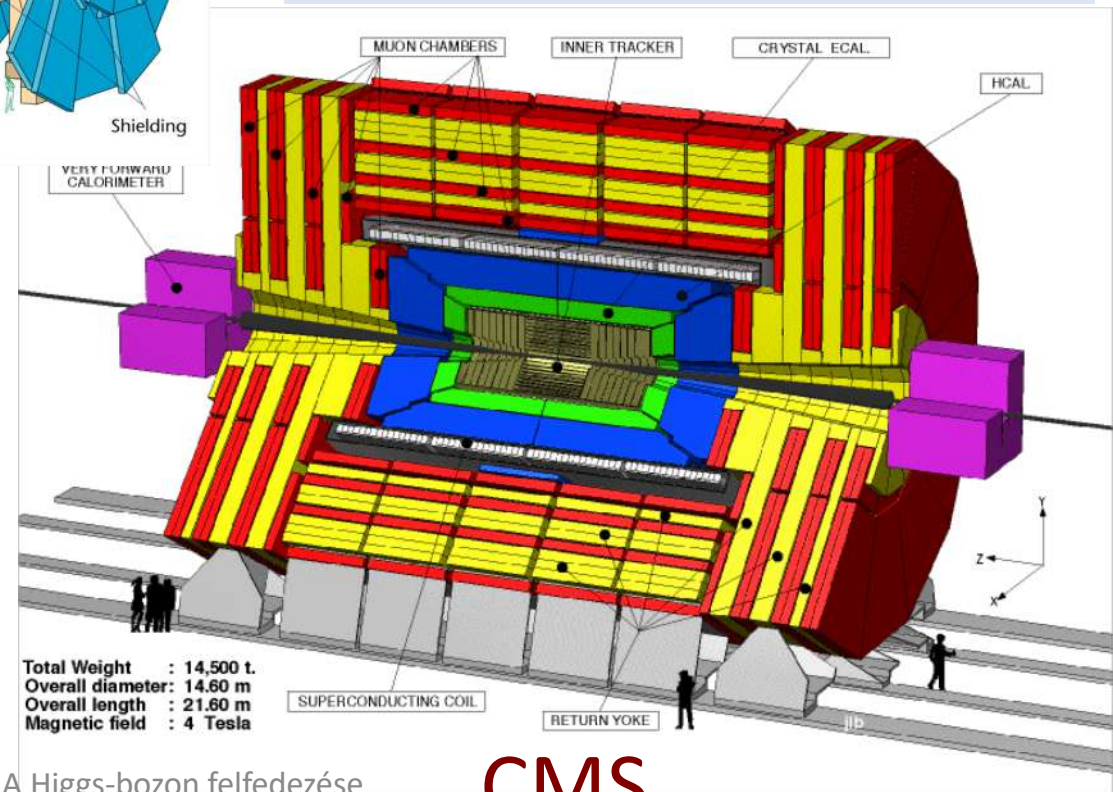
Hungarian technical contribution

CMS:

- Alignment system of muon chambers
- Pixel detector reconstruction, calibration
- Wigner FK Grid Tier2
- Hadron calorimeter (VFCAL) assembly
- High-level trigger
- Luminosity detector calibration, monitoring
- Zero Degree Calorimeter calibration

ATLAS:

- Electromagnetic calorimeter simulation
- Electron-photon reconstruction, trigger

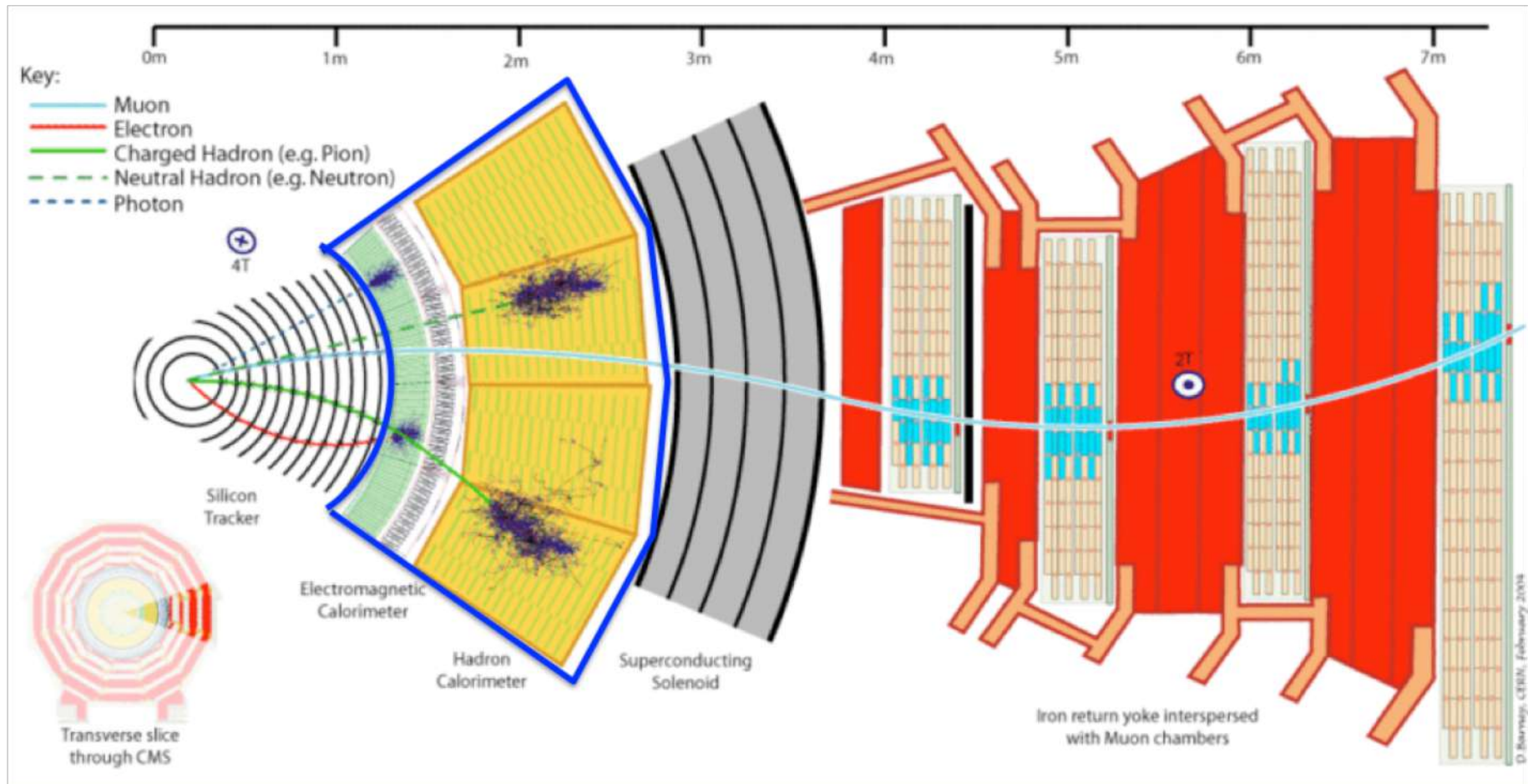


Total Weight : 14,500 t.
Overall diameter: 14.60 m
Overall length : 21.60 m
Magnetic field : 4 Tesla

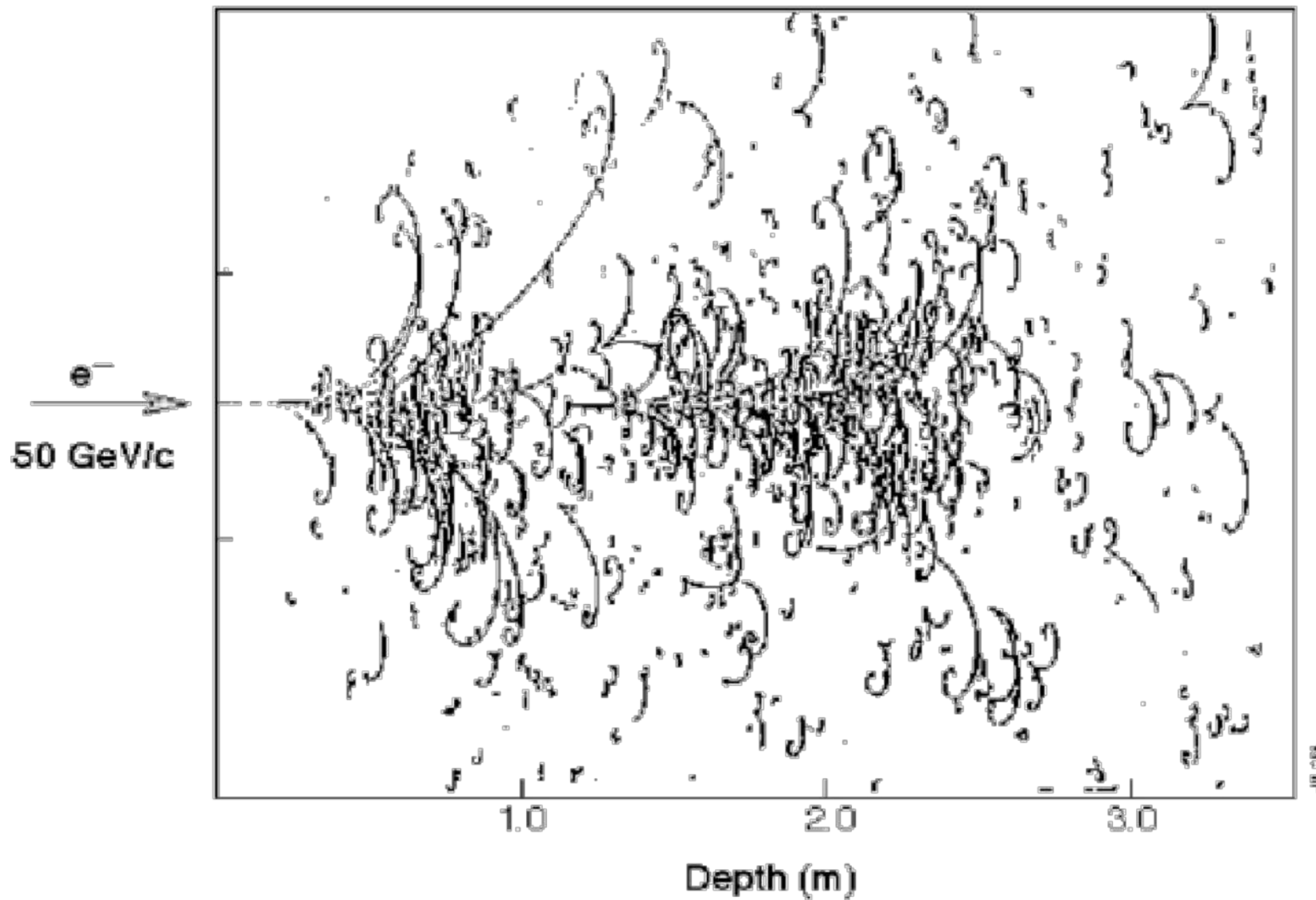
CMS

Calorimeters

- Stop particle (absorb its full particle energy), measure the absorbed energy
 - Thick, high density material placed to the particles' path
- Destructive for almost all particles (except μ , ν) \rightarrow helps to identify muons
- Indispensable to detect neutral particles (photons, neutral hadrons)
- Incoming charged or neutral particles interact electromagnetically with the detector material or create hadronic showers
- Secondary particles ionise or produce excited states in the active material, giving measurable signal



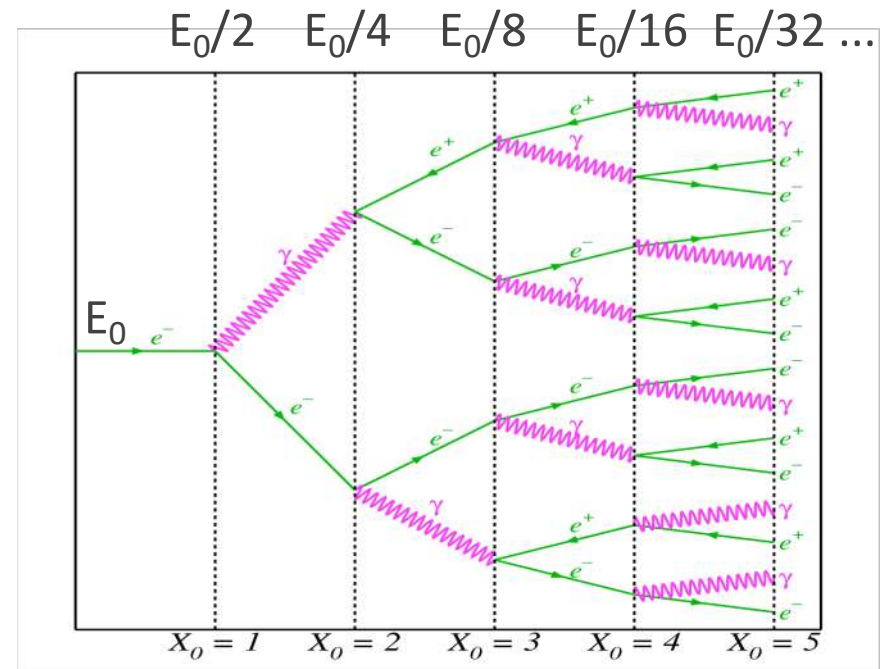
Electromagnetic shower



Homogeneous calorimetry – illustration of bremsstrahlung and photon conversion mechanism:
Electromagnetic shower created by a 50 GeV electron in 3 T magnetic field in the BEBC
bubble chamber using Ne/H₂ (70%-30%) gas mixture

Electromagnetic calorimeters

- $E \gtrsim 100$ MeV:
 - Electron: bremsstrahlung
 $-(dE/dx)_{\text{rad}} = E/X_0$ (X_0 : radiation length)
 energy of e^\pm drops to E/e in X_0 distance
 - Photon: e^+e^- pair productions
 $-(dI/dx) = e^{-x/\lambda} / \lambda$ ($\lambda = 9 X_0 / 7$)
 λ : photon mean free path for pair production
 - EM shower develops until $E > E_{\text{critical}}$
- At lower ($E < E_{\text{critical}} \sim \text{MeV}$) energy, the absorption processes dominate
 - Electron (charged particle): ionisation and excitation of atomic states
 - Photon: photoelectric effect and Compton



Simple shower model:

X_0 as generation length

Distance: $t = x / X_0$

Particle multiplicity: $N(t) = 2^t$

Energy / particle:

$$E(t) = E_0 / 2^t$$

Shower maximum:

$$t_{\text{max}} = \ln(E_0/E_c) / \ln 2$$

Example:

Into a CsI crystal ($E_c \sim 10$ MeV) a 1 GeV electron enters

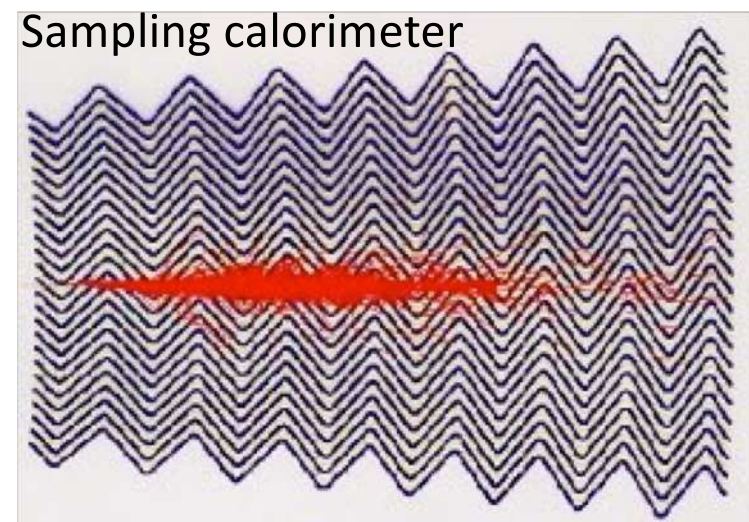
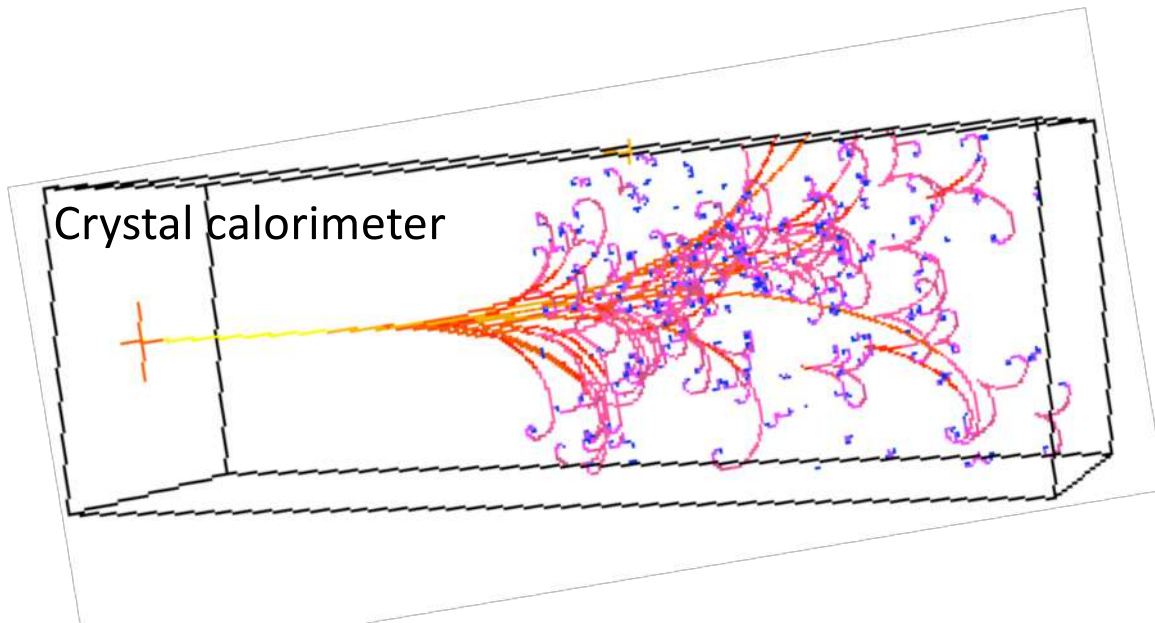
Shower maximum: $N_{\text{max}} = E_0/E_c = 100$, $t_{\text{max}} = 6.6$

Electron with E_c energy fully absorbed in $\sim 1 X_0$ distance but $\sim 7-9 X_0$ material is needed to absorb 95% of the energy of such a photon

→ Calorimeter depth $> 15 X_0$ to cover the full shower

Electromagnetic calorimeters

- Most important properties:
 - Minimum $\sim 15 X_0$ thickness (typically $\sim 25 X_0$)
 - Position of shower maximum depends slowly on energy, calorimeter depth depends logarithmically on energy
 - Energy leakage mostly due to low energy photons escaping the calorimeter cell on the sides (lateral leakage) or behind it (rear / longitudinal leakage)
- In reality shower modelling is much more complicated, modelled by Monte Carlo simulation (ex. GEANT4 program package)



Calorimeter types and processes

- **Homogeneous** (EM) calorimeters: full volume sensitive (contributes to the signal)
 - Usually inorganic (high density, high Z) crystal scintillators or non-scintillator Cherenkov-radiators (e.g. lead-glass)
 - Excellent energy resolution, but expensive and high quality mass-production challenging
 - E.g. CMS lead-tungstate EM calorimeter:

PbWO ₄ (PWO) (CMS)	25X ₀	3%/√E ⊕ 0.5% ⊕ 0.2/E	1997
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- **Sampling** (EM or hadron) calorimeters: absorber metal layers, pl. Fe, Cu (hadron) and Pb (EM), alternate with active detector layers

- Active material could be scintillator, ionizing liquid noble gas, Cherenkov-radiator, semiconductor or gas detector
- Cheap, provides fine segmentation (better vertex determination precision, more e/γ identification information), but not whole shower visible (worse energy resolution)
- E.g. ATLAS EM calorimeter:

Liquid Ar/Pb accordion	25X ₀	10%/√E ⊕ 0.4% ⊕ 0.3/E	1996
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- **Energy resolution:** $(\sigma_E / E)^2 = (a/E)^2 + (b/\sqrt{E})^2 + c$
 - **Fluctuation of first interaction point** a: noise term, c: constant term (→ Calo quality)
 - **Sampling fluctuations**
 - **Fluctuation of energy leakage in longitudinal direction**
 - **crystal non-uniformity**
 - Electronic noise
 - ...
- **Spatial resolution**

Energy resolution

Ideally, if all shower particles counted:

$$E \sim N, \quad \sigma \sim \sqrt{N} \sim \sqrt{E}$$

In practice:

absolute $\sigma = a \sqrt{E} \oplus b E \oplus c$

relative $\sigma / E = a / \sqrt{E} \oplus b \oplus c / E$

a: stochastic term

- intrinsic statistical shower fluctuations

- sampling fluctuations

- signal quantum fluctuations (e.g. photo-electron statistics)

b: constant term

- inhomogeneities (hardware or calibration)

- imperfections in calorimeter construction (dimensional variations, etc.)

- non-linearity of readout electronics

- fluctuations in longitudinal energy containment (leakage can also be $\sim E^{-1/4}$)

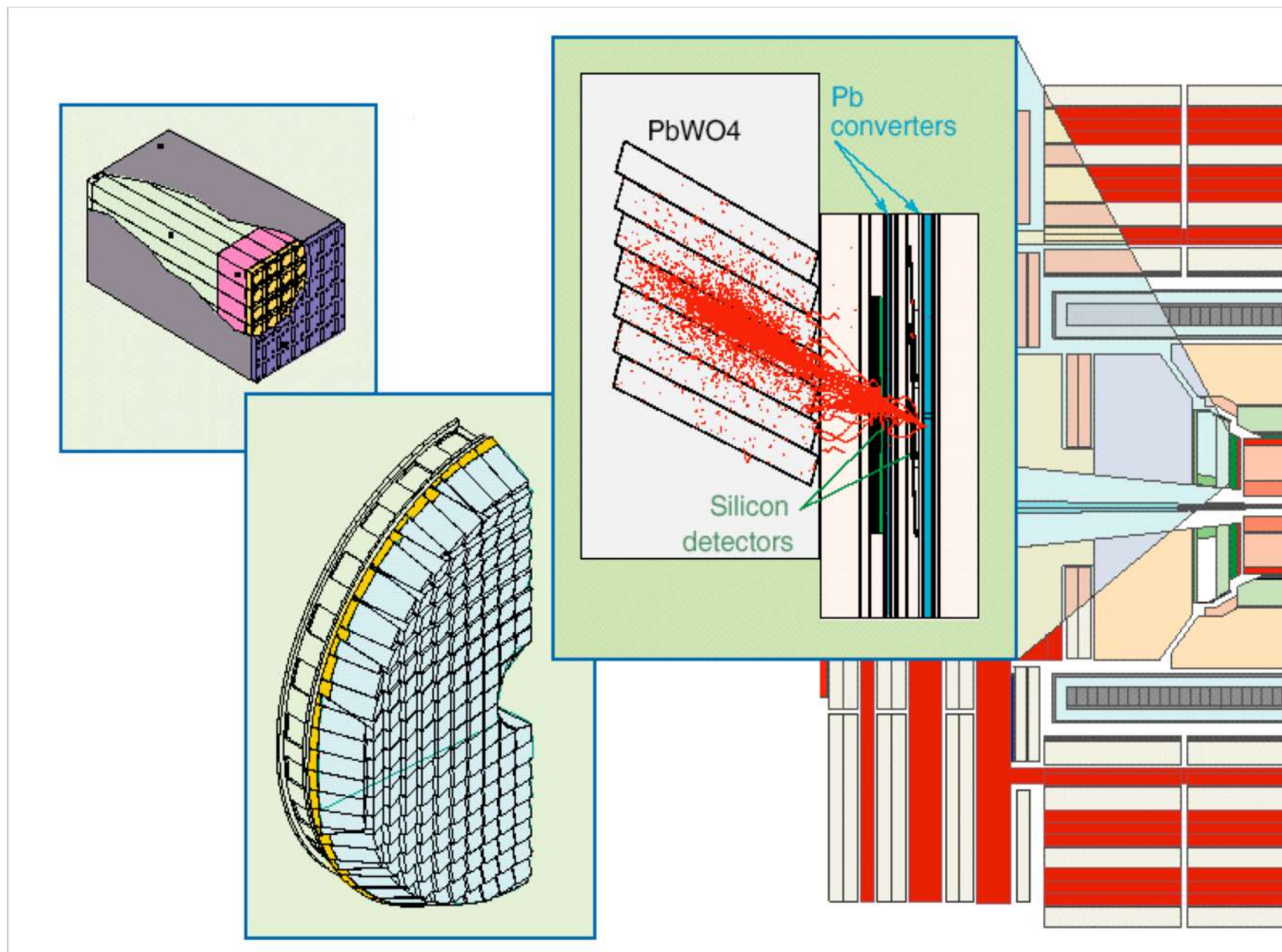
- fluctuations in energy lost in dead material before or within the calorimeter

c: noise term

- readout electronic noise

- Radio-activity, pile-up fluctuations

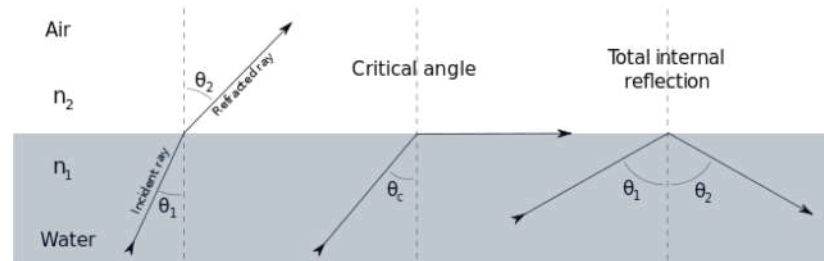
CMS ECAL



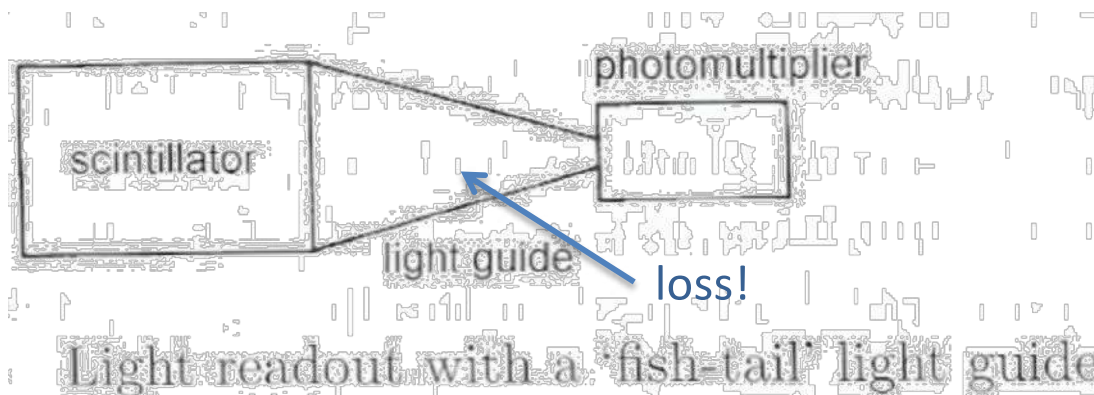
G. Pásztor: Válogott fejezetek a részecskefizikából

Light collection

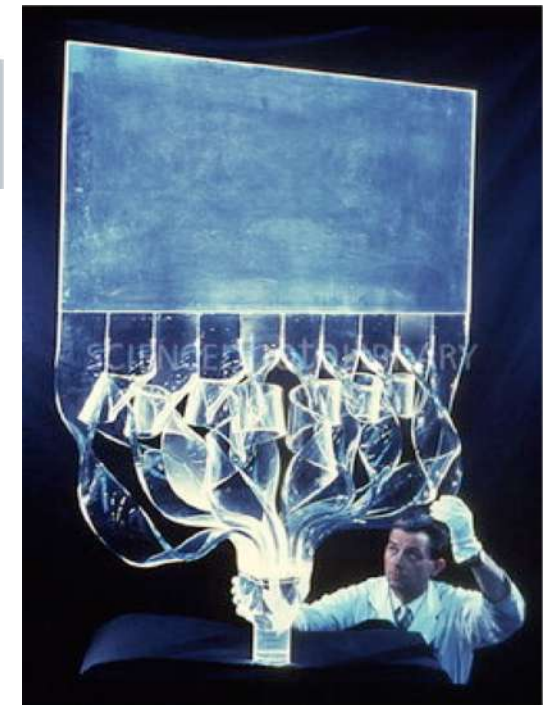
- Usually scintillation counters read out through one or two surfaces
 - Area significantly smaller than the full surface of the scintillator
- For small detectors, the rest of the surface is coated with a reflective diffuse layer (MgO powder or porous Teflon film)
- For large detectors (especially for long rods) total reflection gives the best result (finely polished smooth surface!)



- Plastic scintillators are usually produced in plates
 - The light passing through the edges need to be guided to the photo-detector, fitting there shapes \rightarrow light guides
 - The cross-section can not be reduced without light loss when matched to the photocathode



Light readout with a 'fish-tail' light guide



Adiabatic lightguide
(no light loss)
Only small bending
allowed

Photoelectron multiplier (PM tube)

- Measure fast light signal: light \rightarrow electric signal
- Visible or UV light frees electrons from photocathode via photoelectric effect
- Usually semi-transparent photocathode: very thin semiconductor (SbCs, SbKCs, SbRbKCs) layer on the inner surface of the transparent input window
- The large kinetic energy electron kicks out several others when hitting the dynode ($<25 e / 200 eV$)

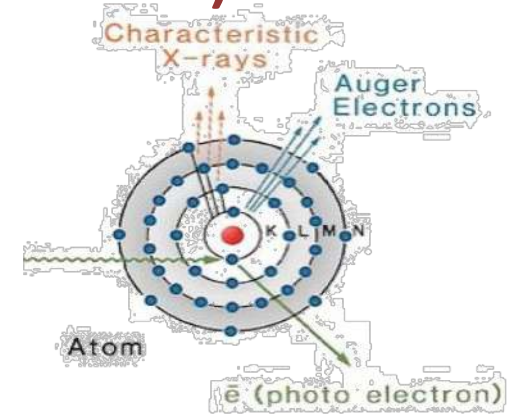
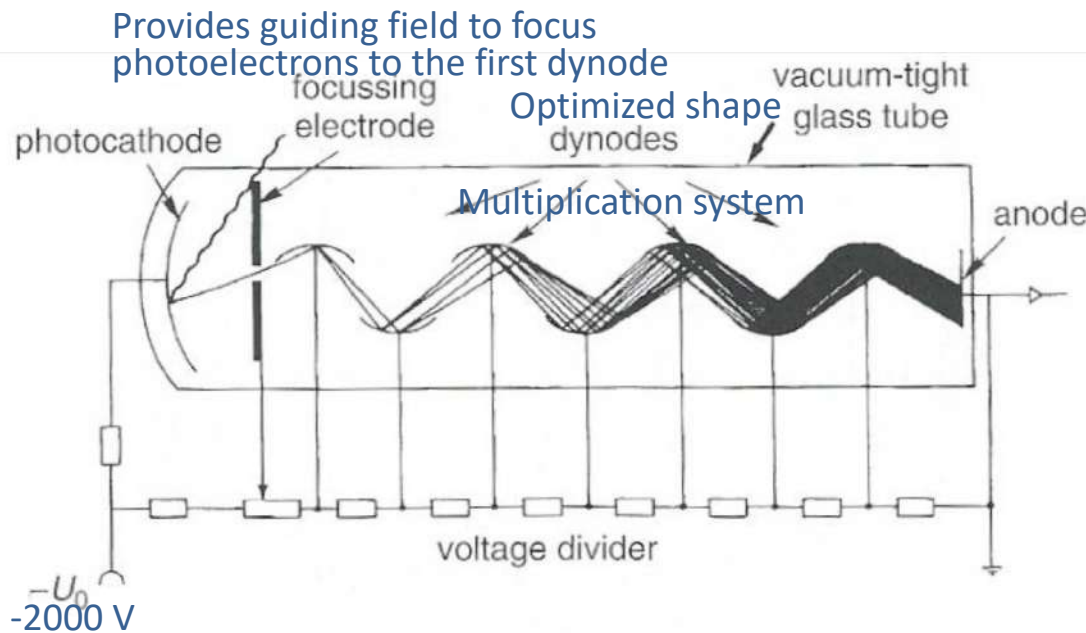


Illustration of the photoelectric effect.

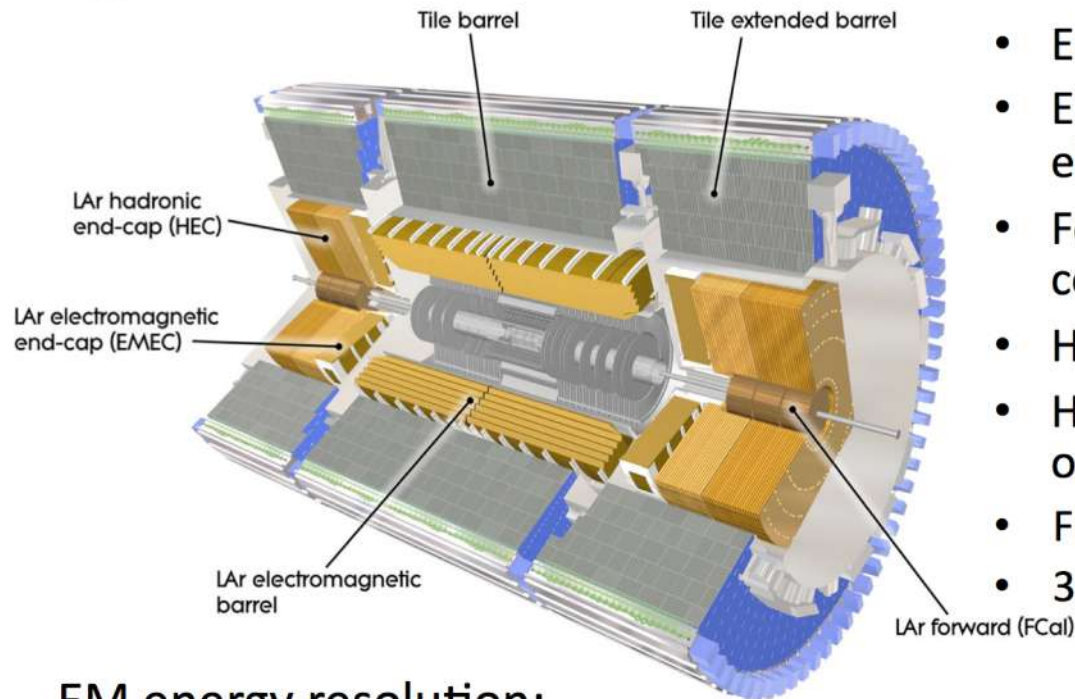


- Quantum efficiency – average number of photoelectrons produced by an incoming photon: $\sim 25\%$, max. 50% (GaAs, GaInAsP)
- Secondary emission coefficient ($g \sim 4$) \rightarrow current amplification: $A = g^n$ (n dynodes)
variance: $(\sigma/A)^2 = 1/(g-1)$
- Total amplification: $>10^6$
- Total transit time: $\sim 10-40$ ns
Rise time: $1-3$ ns

Fig. 5.29. Working principle of a photomultiplier. The electrode system is mounted in an evacuated glass tube. The photomultiplier is usually shielded by a mu-metal cylinder made from high-permeability material against stray magnetic fields (e.g. the magnetic field of the Earth).



The ATLAS Calorimeter System



- EM presampler, $|\eta| < 1.8$: active LAr layer
- EM calorimeter, $|\eta| < 3.2$: accordion shaped electrodes and lead absorbers in LAr
- Forward calorimeter, $3.1 < |\eta| < 4.9$: copper/tungsten-LAr
- Hadronic end-cap, $1.5 < |\eta| < 3.2$: copper/LAr
- Hadronic tile calorimeter, $|\eta| < 1.7$: outside LAr cryostats
- Fine lateral segmentation
- 3 or 4 longitudinal layers

EM energy resolution:

$$\frac{\sigma_E}{E} = \frac{a}{\sqrt{E}} \oplus \frac{b}{E} \oplus c$$

$a \approx 10\%$, $b \approx 170 \text{ MeV}$, $c \approx 0.7\%$

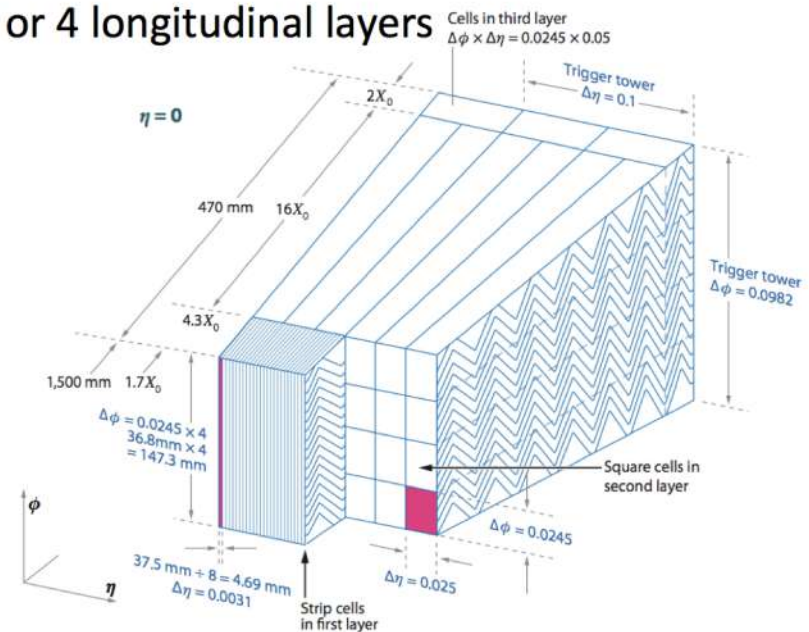
Hadronic energy resolution:

$a \approx 50\%$, $c \approx 3\%$

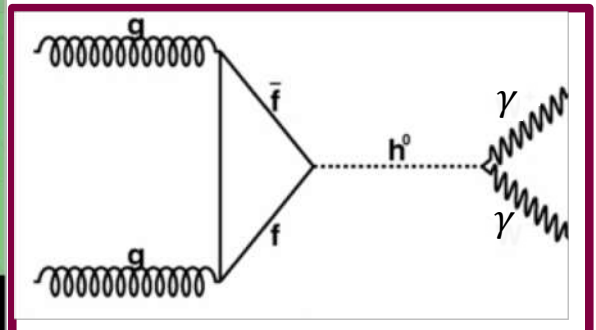
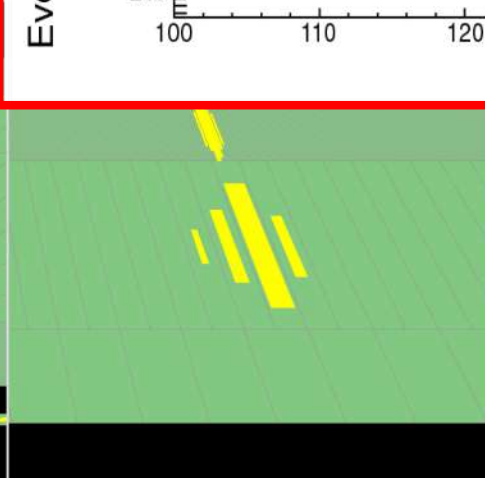
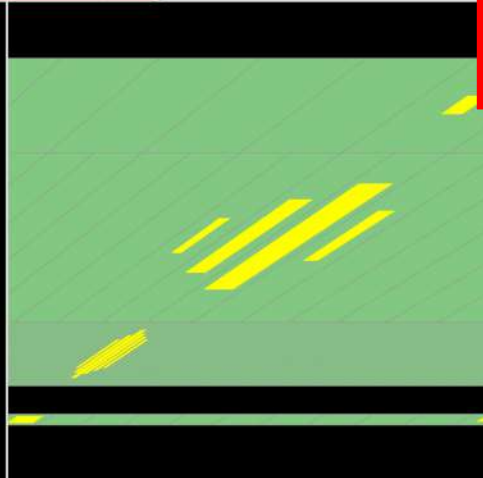
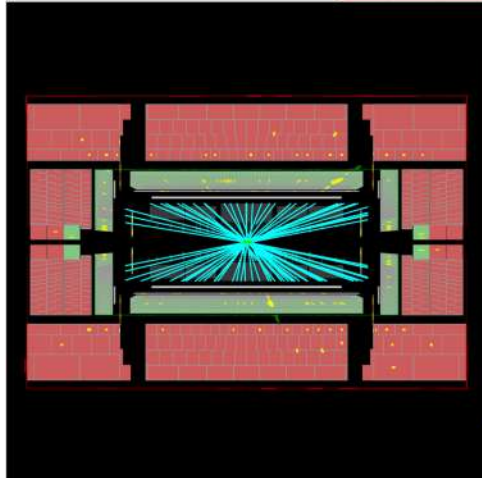
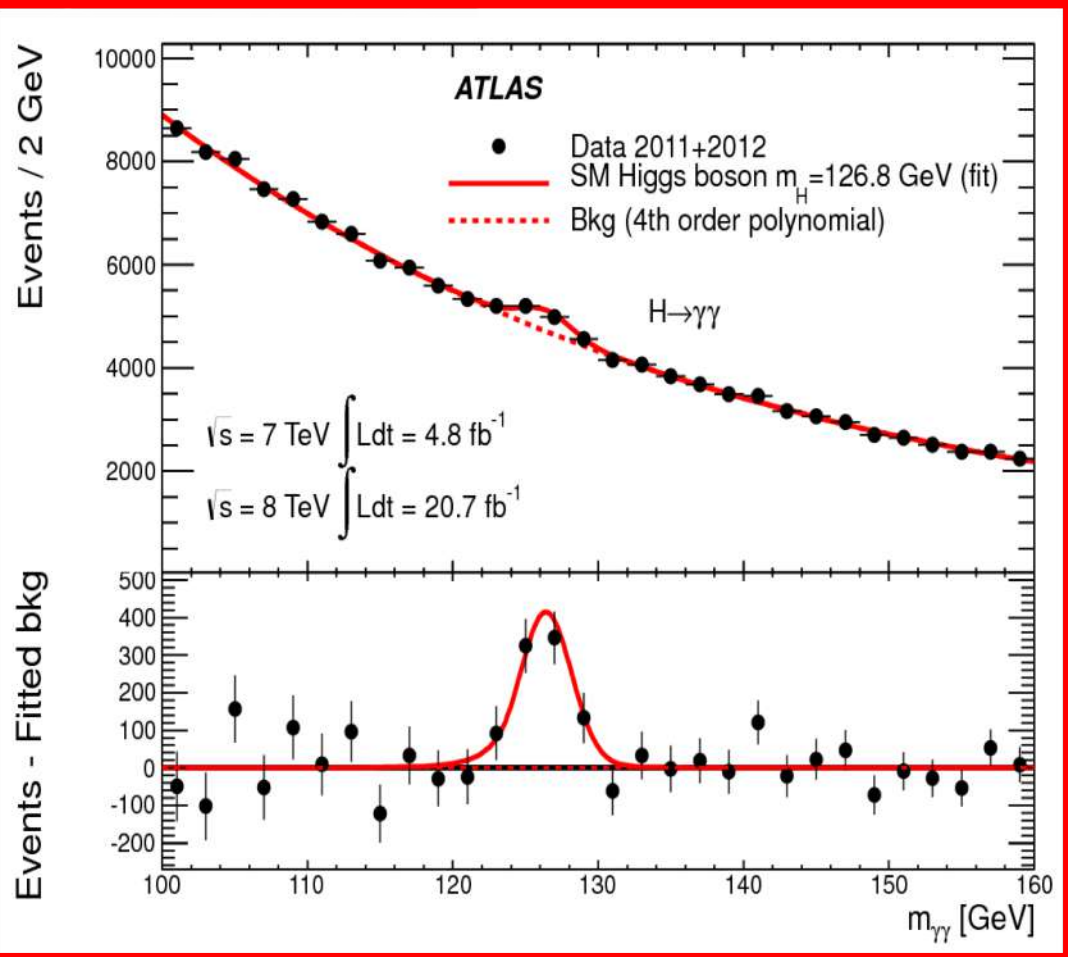
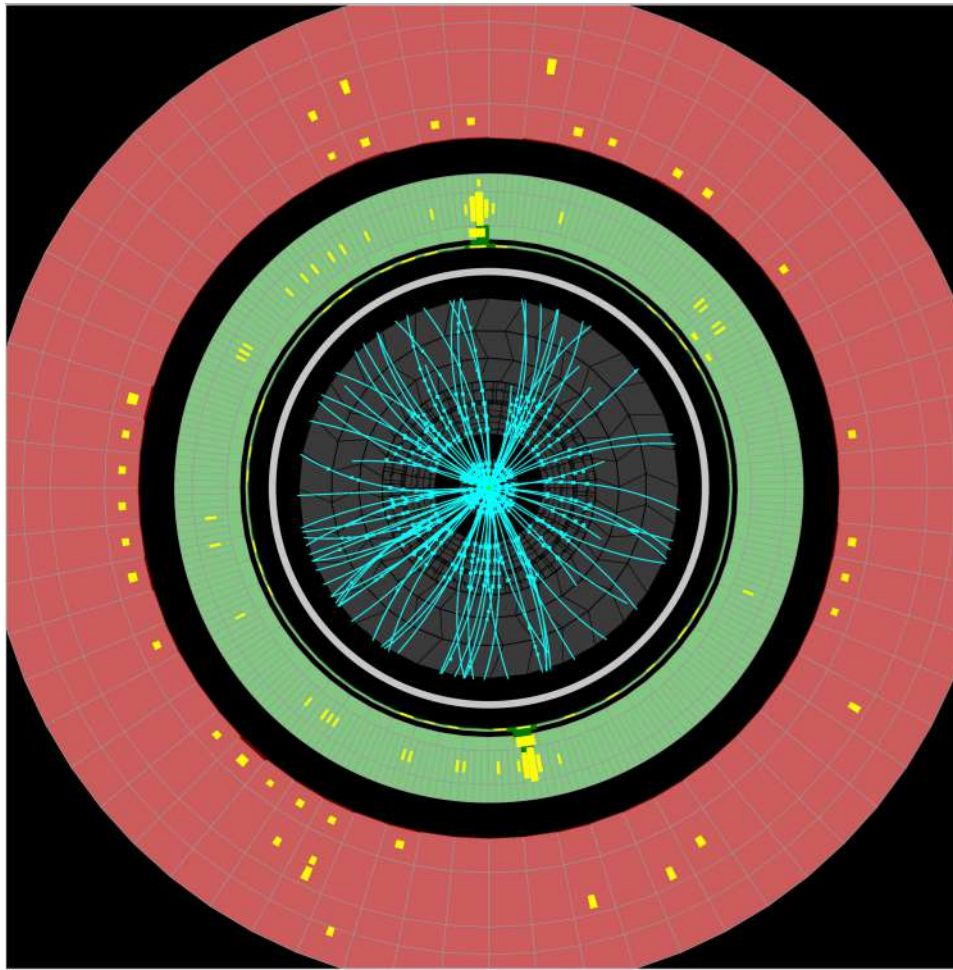
Compare with CMS crystal size:

22 x 22 x 230 mm³ (1 R_M x 1 R_M x 26 X_0) (barrel),

30 x 30 x 220 (endcap)



H \rightarrow $\gamma\gamma$

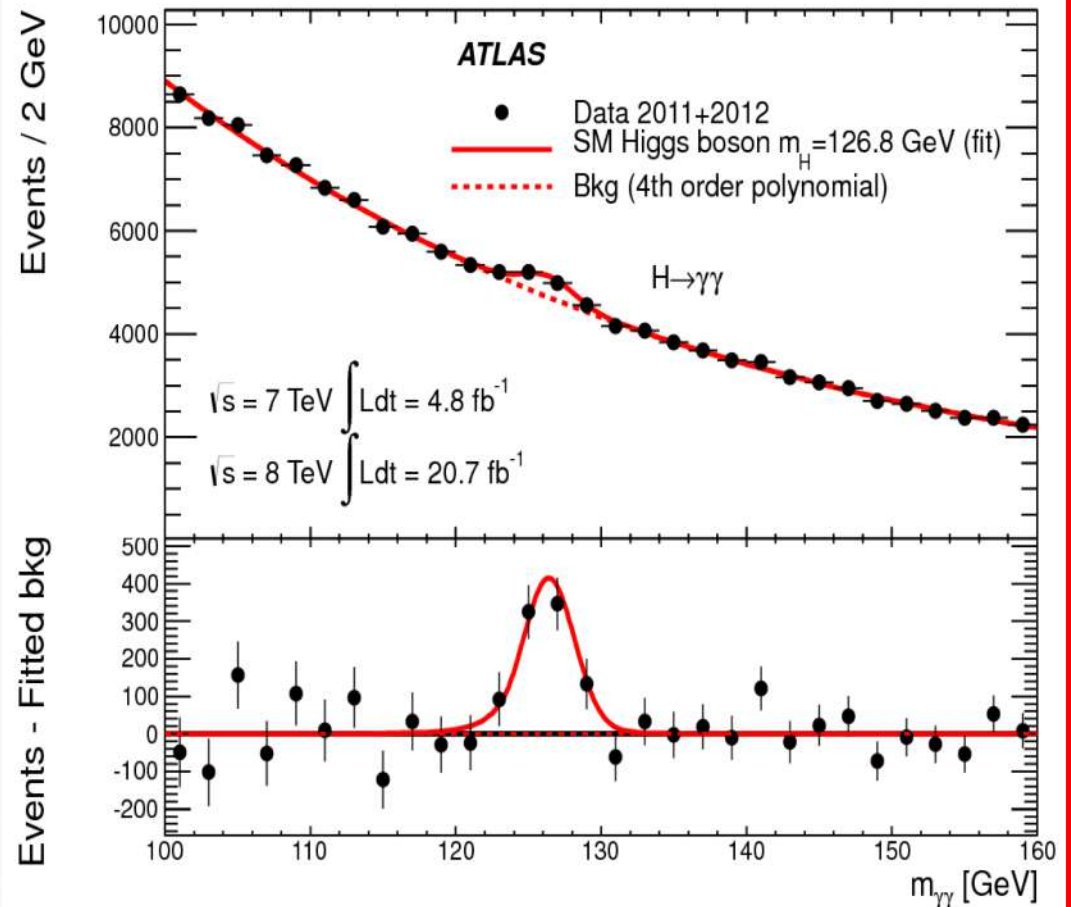


H → γγ

Precise mass reconstruction
 $\sigma_m/m = 0.5 [\sigma_{E1}/E_1 \otimes \sigma_{E1}/E_1 \otimes \cot(\theta/2)\Delta\theta]$
 requires excellent energy resolution and
 vertex position determination!

CMS: excellent energy resolution with
 homogenous crystal calorimeter

ATLAS: photon direction measurement
 with finely segmented calorimeter



Homework

- 100 GeV electron enters (a) a Copper or (b) an Iron absorber. Where will be the shower maximum? How many particles will form the shower at its maximum? What is the average energy of the particles there?
- Compare the electromagnetic calorimeters of the ATLAS and CMS experiments. What technologies are used? How the energy and spatial resolutions compare? What do you think the arguments were when justifying the experiment's own designs?
- EM calorimeter has a stochastic term to its energy resolution of $0.05/\sqrt{E}$. How can we ensure that the energy resolution for a $E = 40$ GeV photon does not exceed 1%?
- How will interact a muon with the detector material when it traverses through the CMS detector? How can we use the signals to identify the muon? Do the calorimeters play any role?
- A multi-purpose detector system using the traditional layered structure has a relative track momentum resolution of $0.00015 \text{ pT [GeV]} \oplus 0.005$, an EM calorimeter energy resolution of $0.2/E[\text{GeV}] \oplus 0.03/\sqrt{E[\text{GeV}]} \oplus 0.005$ and a hadron calorimeter energy resolutions of $0.7/\sqrt{E[\text{GeV}]} \oplus 0.08$. At what electron and pion energy will the tracking and the calorimeter measurements have the same precision for highly relativistic particles at pseudorapidity of 0? Which measurement is more precise at low / high momentum?